



Research Article

VIBRATIONAL ANALYSIS AND HYDROGEN-BONDING EFFECTS ON THE VIBRATIONAL MODES OF ZWITTERIONIC DL-TRYPTOPHAN: IR SPECTROSCOPY AND DFT CALCULATIONS

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ABSTRACT

The vibrational behavior of zwitterionic DL-tryptophan has been investigated using infra-red spectroscopy and density functional theory calculations. In order to study the hydrogen bonding effects on the molecular structure and vibrational normal modes of the under-investigated amino acid, both the solid state and aqueous solution calculations have been presented. The three dimensional molecular structure optimized by the solid state calculations shows the best agreement with that previously published using X-Ray diffraction technique. In addition, the calculated and experimentally observed wavenumbers have been compared. However, only the calculations performed in the solid state allowed us to attribute the band observed at around 2532 cm⁻¹ to the N-H stretching vibration. Because some normal modes are strongly affected by the extensive hydrogen bonding network present in DL-tryptophan crystal, the solid state model is the most suitable for interpreting the experimental infra-red spectrum.

Keywords: DL-tryptophan, amino acids, IR spectroscopy, DFT calculations.

1. INTRODUCTION

Amino acids are the main building blocks of proteins, they play a very important role in almost all the biological mechanisms in Nature. The molecular structures of amino acids have both carboxyl and amino groups as substituents on the same carbon atom called α -carbon. However, the amino acids with both positive (NH⁺₃) and negative (COO⁻) electrical charges are known as zwitterionic amino acids [1]. With the exception of glycine, amino acids can exist in two enantiomeric forms: L or D-forms. A mixture of equal amounts of the two forms of amino acids is called a racemic compound (DL-form) [2].

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The ability to distinguish between the L, D and DL-forms of amino acids is naturally manifested in our biological world. It has been shown that "almost" only the L-form is naturally present in proteins [3]. In fact, this intimate property of our biological World remains a mystery and its origin is still unknown until now [4,5].

Since the early 20th century, tryptophan has been the subject of intensive spectroscopic studies [6]. Some of them showed that the ability of proteins to absorb the UV radiations is mainly due to the presence of amino acids, and tryptophan has the strongest absorption near 280 nm [7]. The IR spectra of L-tryptophan were reported by several researchers [8,9]. However, the detailed vibrational analysis was performed by X. Cao et al. [10].

From a crystallographic point of view, it has been experimentally confirmed that the crystal structures of DL-tryptophan [11,12] and L-tryptophan [13] are completely different, they crystallize in the monoclinic (space group $P2_1/c$) and triclinic (space group $P1$) systems, respectively. Recent studies have also revealed changes in the spectroscopic behavior of many L- and DL-forms of amino acids, such as: L- and DL-histidine [14], L- and DL-alanine [15] and L- and DL-serine [16,17]. All these studies showed that the differences between the vibrational spectra of L- and DL-amino acids are mainly caused by the peculiarities of their inter-molecular interactions. It is important to note that each amino acid crystal has its specific hydrogen bond (H-bond) network, which plays a major role in determining the most stable molecular conformation. Theoretically, the inclusion of solvation effects in the calculations for predicting the most stable molecular conformations of amino acids have been widely employed. However, such approaches can encounter two difficulties: (i) the first one, is the large number of possible conformers which characterizes the amino acid molecules [18,19], (ii) the second difficulty is that the effects of inter-molecular interactions which appear, in general, in the region higher than $\sim 1600\text{ cm}^{-1}$ in the experimental IR spectra of amino acids (KBr pellet techniques) [10,20,21] cannot be correctly taken into account in the calculations.

Experimentally, several methods based on the matrix isolated monomeric molecules of amino acids have been developed to improve the agreements between the observed and calculated IR spectra [10]. However, such methods eliminate the "real effects" of inter-molecular interactions which can be differentiated the IR spectra of L, D- and DL-amino acids.

In our study, the IR spectrum recorded in the solid phase has been studied to reveal the effects of inter-molecular interactions on the vibrational modes of the considered compound, and the periodic (solid state) calculations have been employed to overcome the problems of molecular conformations and inter-molecular interactions.

The main aim of this work is to better understand the effects of hydrogen bonding on the structural and vibrational normal modes of DL-tryptophan. To this end, its molecular structure has been predicted in both solution and the solid state and compared with previously reported experimental results. Also, for the first time, the IR spectrum of DL-tryptophan, recorded in the region $600\text{--}4000\text{ cm}^{-1}$, has been fully studied employing non-periodic and periodic calculations.

2. MATERIALS AND METHODS

2.1. Experimental details

The zwitterionic form of DL-tryptophan ($\geq 99\%$) was supplied by Sigma Aldrich and utilized without further purification.

At room temperature, the FT-IR spectrum of powder of DL-tryptophan was recorded in the region $600\text{--}4000\text{ cm}^{-1}$ using solid-state sample, on a JASCO FT/IR-6300 spectrometer. The spectral resolution of the spectrometer was 4 cm^{-1} .

2.2. Computational details

For the aqueous-solution model, DFT calculations were carried out with the ORCA 3.0.3 program [22], with the use of water as solvent. The initial geometry of DL-tryptophan was generated without any constraints from single crystal X-ray diffraction results [12]. The optimized geometrical parameters, vibrational wavenumbers and IR intensities were calculated by using the hybrid DFT method B3LYP [23,24], the def2-TZVP basis set [25,26] and also the def2-TZVP/J auxiliary basis set. In these calculations, the concerning energy terms were obtained by using the RIJCOSX algorithm [27,28]. The systematic errors due to the neglect of the anharmonic effects were compensated by using the correlation expression below, which is similar to those proposed by Osaki and Soejima [29] and Yoshida et al. [30].

$$\omega_{\text{scal}} = -0.000028 \omega_{\text{cal}}^2 + 1.034 \omega_{\text{cal}} - 14.77 \quad (\text{cm}^{-1}) \quad (1)$$

Here, ω_{scal} and ω_{cal} are the scaled and calculated wavenumbers for each vibrational mode, respectively.

The solid phase calculations for DL-tryptophan in crystal structure were realized in the GGA (generalized-gradient-approximation) by using the PBE (Perdew–Burke–Ernzerhof) functional [31] in the CASTEP program [32]. In the calculations, the norm-conserving pseudopotentials and the plane-wave cut-off energy of 830 eV were employed. Brillouin zone sampling of electronic states was performed on $1 \times 2 \times 2$ Monkhorst Pack grids. The harmonic vibrational wavenumbers were obtained by using the optimized structure and employing the linear response method within the density functional perturbation theory (DFPT) [33] at the Gamma point.

3. RESULTS AND DISCUSSION

3.1. Geometry optimization

The optimized molecular structure of DL-tryptophan in aqueous solution is shown in Figure 1a. Figures 1b and 1c represent the optimized unit cell and the molecular packing restricted by inter-molecular N-H \cdots O H-bonds in the solid phase, respectively. The observed (X-Ray) [12] and computed unit-cell parameters are listed in Table I.

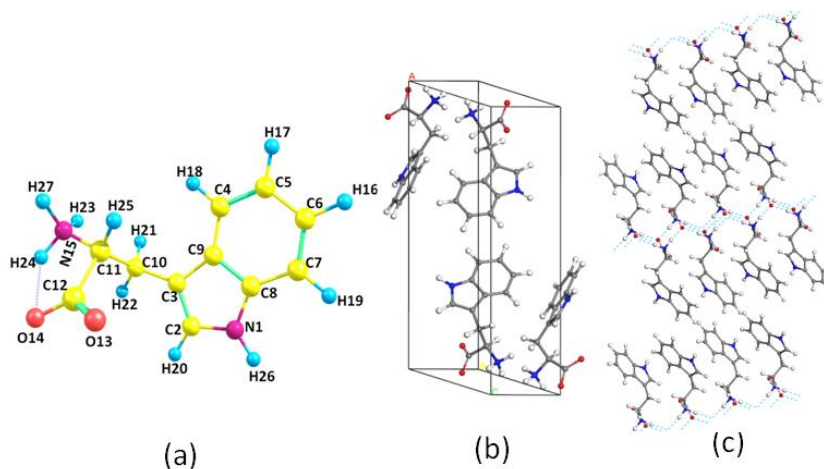


Figure 1. The optimized geometries obtained at the Density Functional Theory approach for the molecular and crystal structures of DL-tryptophan. (a) the optimized molecular structure obtained by the aqueous solution model; (b) the optimized geometry obtained for the unit cell of DL-tryptophan by using the solid phase model; (c) the molecular packing obtained by using the solid phase model.

Table 1. A comparison between the calculated and experimentally measured unit cell parameters obtained for DL-tryptophan crystal.

| Unit cell parameters | Calculated (a) | Experimental (b) |
|----------------------|----------------|------------------|
| a (Å) | 18.9602 | 18.899(2) |
| b (Å) | 5.7709 | 5.7445(6) |
| c (Å) | 9.1280 | 9.309(1) |
| β (°) | 100.956 | 101.776(2) |
| V (Å ³) | 980.56 | 989.4(4) |

(a) The modelled unit cell is shown in Figure 1b.

(b) The experimental unit cell parameters were taken from ref. [12].

Our results show that the calculated unit-cell parameters are in good agreement with the ones obtained by using X-ray diffraction ; the unit-cell volume calculated with the PBE method is about 0.9 % less than the experimentally observed one.

As can be seen in Figure 1c, the NH_3^+ group forms three intermolecular $\text{N-H}\cdots\text{O}$ H-bonds with carboxyl oxygen atoms of adjacent molecules, the computed distances between donor and acceptor atoms are ranging from 2.684 to 2.821Å. It is important to mention that the stability of DL-tryptophan in the solid state is strongly dependent on this extensive network of H-bonds. Therefore, significant differences between the computed dihedral angles in aqueous solution and the corresponding ones obtained experimentally can be expected.

Some selected optimized geometry parameters of DL-tryptophan molecule (bond lengths, bond angles and dihedral angles) are listed in Table II in comparison with the corresponding experimental parameters [12]. The numbering scheme used for the atoms in the table is as given in Figure 1a.

Table 2. A comparison between the calculated and experimentally measured geometry parameters obtained for DL-tryptophan molecule.

| Bond lengths (Å) | | | |
|----------------------------|--------------------------|-----------------------|-------------------|
| Parameters (a) | Solution phase model (b) | Solid phase model (c) | Experimental [12] |
| N1C2 | 1.373 | 1.372 | 1.373(2) |
| C2C3 | 1.3701 | 1.376 | 1.359(2) |
| C3C9 | 1.4393 | 1.439 | 1.435(2) |
| C3C10 | 1.4957 | 1.484 | 1.490(1) |
| C4C5 | 1.3861 | 1.393 | 1.375(2) |
| C4C9 | 1.4034 | 1.402 | 1.393(2) |
| C5C6 | 1.4076 | 1.410 | 1.397(2) |
| C6C7 | 1.3866 | 1.390 | 1.374(2) |
| C7C8 | 1.3955 | 1.399 | 1.391(2) |
| C8C9 | 1.4213 | 1.428 | 1.414(1) |
| C10C11 | 1.5428 | 1.546 | 1.537(1) |
| C11N15 | 1.5042 | 1.456 | 1.486(1) |
| C11C12 | 1.5544 | 1.534 | 1.522(1) |
| C12O14 | 1.2601 | 1.270 | 1.263(1) |
| C12O13 | 1.2423 | 1.2 | 1.238(1) |
| Bond angles (°) | | | |
| C3 C3C10C11 | 115.048 | 114.84 | 113.71(8) |
| N15C11C12 | 106.370 | 110.12 | 109.50(8) |
| N15C11C10 | 108.352 | 109.84 | 109.61(8) |
| C12C11C10 | 112.521 | 111.14 | 112.16(8) |
| O13C12O14 | 128.215 | 126.11 | 125.95(9) |
| O13C12C11 | 116.738 | 119.25 | 117.60(8) |
| O14C12C11 | 115.033 | 114.63 | 116.45(9) |
| Dihedral angles (°) | | | |
| C2C3C10C11 | 107.59 | 104.00 | 107.9(1) |
| C3C10C11N15 | 171.88 | 163.3 | 166.78(8) |
| C3C10C11C12 | -70.79 | -74.57 | -71.4(1) |
| N15C11C12O13 | -155.30 | -157.90 | -158.0(1) |
| N15C11C12O14 | 25.95 | 22.29 | 21.7(1) |
| C10C11C12O14 | -92.53 | -99.65 | -100.3(1) |

(a) The atom numbering scheme is given in Figure 1a.

(b) The calculation model used for the aqueous solution phase of DL-tryptophan.

(c) The calculation model used for DL-tryptophan in solid phase.

It can be seen from Table II that the computed bond lengths for both solution and solid phase models are almost in good agreement with the observed results, the main deviation between them is found to be 0.03 Å for the C11-C12 bond length.

The same observation holds true for the bond angles, where the root mean square (rms) deviation between the predicted and observed bond angles are about 1.74° (solution model) and 1.12° (solid phase model).

In contrast, the calculated C10-C11-C12-O14 dihedral angle at -92.53° in aqueous solution and experimentally observed at -100.3(1)° is the most affected by the H-bonding network. This difference (~7.77°) is due to the effect of the intra-molecular N15H24... O14 H-bond (Figure 1a), noting that this H-bond is missing in the crystal of DL-tryptophan [12]. In the solid state model, this difference is reduced to 0.6°. On the other hand, the differences found between the computed (aqueous solution model) and the observed N15-C11-C12-O13 and N15-C11-C12-

O14 dihedral angles are 2.7° and 3.6° , respectively. These differences are attributed, as expected above, to the effects of inter-molecular (H-bonds) interactions which are not correctly taken into account in the solution-phase calculations.

With only a few exceptions, such as the C2-C-3C10-C11 and C3-C10-C11-C12 dihedral angles, it can be concluded that the computed geometrical parameters using the solid state calculations show the best agreement with the observed molecular structure.

3.2. Vibrational analysis

In this section, in which the assignments of the vibrational bands observed in the FT-IR spectrum of DL-tryptophan are presented and evaluated, a comparison of the experimental results with those previously reported for L-tryptophan is given and the simulated IR spectra for the title compound in solution and in solid phase circumstances are presented; please see Figure 2 for the recorded experimental FT-IR spectrum of DL-tryptophan.

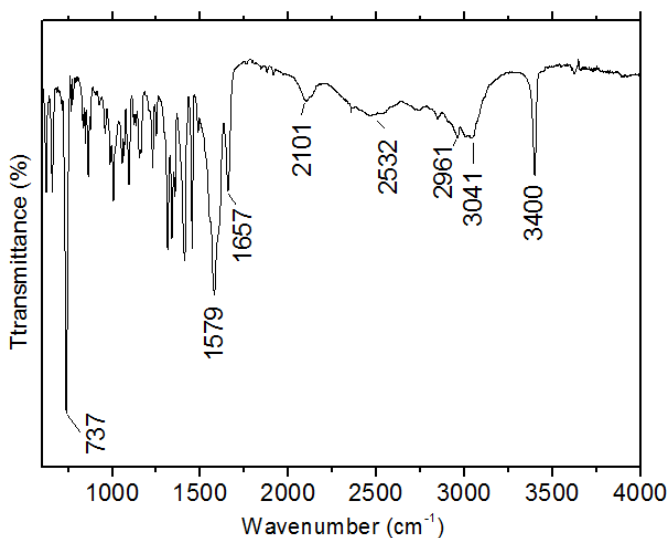


Figure 2. The experimental IR spectrum of DL-tryptophan

As a first observation, the FT-IR spectra recorded for the DL-tryptophan (Figure 2) and L-tryptophan (KBr disc technique)[10] are similar to each other, but not exactly identical. As the band intensities are considered, it is seen that the strongest band in the FT-IR spectrum of DL-tryptophan is centered at 737 cm^{-1} , while the strongest three bands observed in the IR spectrum of L-tryptophan, which have almost the same intensity, are centered at 744 , 1590 and 3402 cm^{-1} .

It is important to note that there is a general tendency of down-shifting (or red-shifting) for the positions of the bands observed in the FT-IR spectrum of DL-tryptophan in comparison to the corresponding IR bands observed for L-tryptophan. However, the IR band observed at 2100 cm^{-1} (for L-tryptophan, the IR band at 2073 cm^{-1}) is an exception for this situation.

Moreover, it is remarkable that the region between 1657 and 2845 cm^{-1} in the observed FT-IR spectrum of DL-tryptophan, is characterised by some broad-bands. This region is similar to those observed for the majority of IR spectra of amino acids, they almost all have a particular band at around 2100 cm^{-1} [10,19-21,34]. However, an assertive interpretation of the origin of this band has not yet been given, for example; the observed band at 2040 cm^{-1} in the IR spectrum of

L-serine was assigned to the combined vibration δNH_3^+ and τNH_3^+ [16], the band observed at 1991 cm^{-1} for L-histidine perchlorate [35] was attributed to the overtones and combinational tones, moreover, a similar band was appeared at 2073 cm^{-1} for L-tryptophan (KBr pellet method), but it was not observed in the spectrum recorded by using the dissolution-spray-deposition (DSD) technique [10].

In the following, the calculated IR spectra using both the solid state and aqueous solution models will be presented and analyzed to get better understanding of the vibrational behavior of DL-tryptophan.

The DL-tryptophan molecule belongs to C_1 symmetry point group, it has 27 atoms, which gives $3n - 6 = 75$ normal modes of vibration (n is the number of atoms in the molecule). The aqueous solution computations give the harmonic vibrational wavenumbers of the 75 internal modes as they are given in Appendix A. In order to obtain a good agreement between the calculated and observed wavenumbers, the correlation relation given by the equation (1) was used. Figure 3 shows the quadratic regression plot between the experimental and theoretical wavenumbers which is fitted into a quadratic equation. The theoretical IR spectrum obtained in the aqueous solution model, which is shown in Figure 4a, was produced over the harmonic wavenumbers fitted by using this correlation relation.

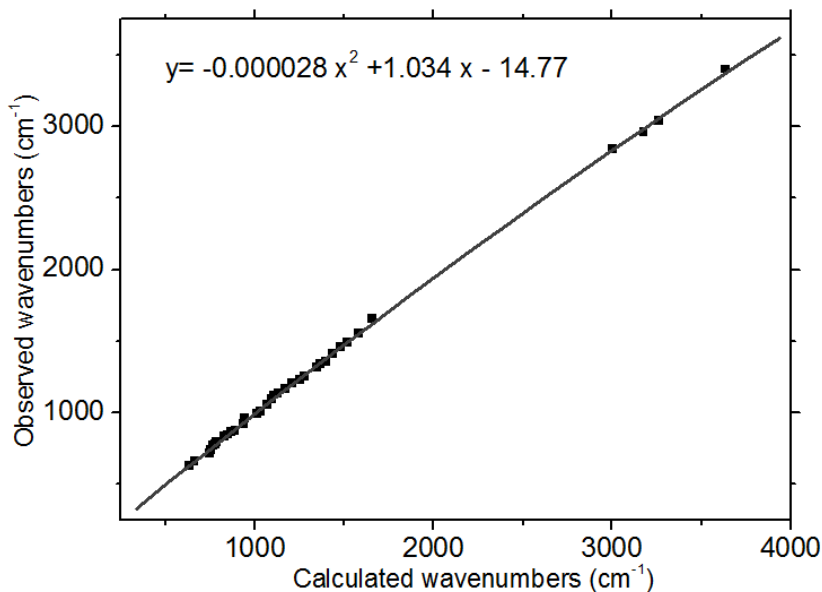


Figure 3. The correlation between the calculated and observed wavenumbers for DL-tryptophan molecule.

For the solid phase calculations, the unit cell contains four molecules (108 atoms), each atom has three degrees of freedom, therefore there are 324 modes distributed between internal modes, lattice modes, librational modes and three acoustic modes. The acoustic modes are computed as small values close to zero at -0.063 , -0.051 and -0.037 cm^{-1} , the remaining 321 optical modes are classified, according to the irreducible representation of the C_{2h} factor group, into four subgroups a_g , b_g , a_u and b_u as can be seen in Appendix B. Where: a_u and b_u are the IR active modes, and a_g and b_g are the RAMAN active modes [36]. In our case, the irreducible representations of the optical modes are given as: $\Gamma_{Optical} = 81a_g + 81b_g + 80a_u + 79b_u$.

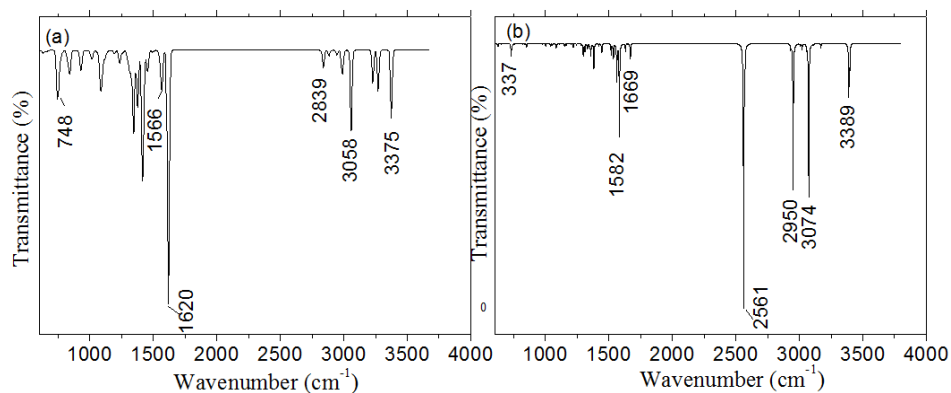


Figure 4. The infrared (IR) spectra calculated for DL-tryptophan by using the Density functional theory. (a) the scaled theoretical IR spectrum obtained by the aqueous solution model; (b) the theoretical IR spectrum obtained by the solid-state model.

The observed and calculated wavenumbers in the region 600-4000 cm^{-1} with their assignments are listed in Table III.

Table 3. A comparison between the experimental and calculated wavenumbers obtained for DL-tryptophan.

| $\omega(\text{obs})$ (a) | Solution model (b) | | | Solid phase model (c) | | | | |
|-----------------------------|-----------------------|-----------------------|---|-----------------------|----------------|-----------------|----------------|---|
| | $\omega(\text{calc})$ | $\omega(\text{scal})$ | Assignment | a_g modes | b_g modes | a_u modes | b_u modes | Assignment |
| 3400 | 3636 (204) | 3375 (204) | $\nu \text{N}_1\text{H}_{26}$ | 3389 (00) | 3396 (00) | 3397 (1163) | 3389 (2039) | $\nu \text{N}_1\text{H}_{26}$ |
| 3041 | 3260 (245) | 3058 (245) | $\nu \text{N}_{15}\text{H}$ | 3045 (00) | 3079 (00) | 3057 (181) | 3074 (6062) | $\nu \text{N}_{15}\text{H}$ |
| 2961 | 3176 (48) | 2987 (48) | νCH | 2933 (00) | 2935 (00) | 2965 (7) | 2950 (5782) | $\nu \text{N}_{15}\text{H}$ |
| 2845 | 3004 (54) | 2839 (54) | νCH_2 | 2930 (00) | 2931 (00) | 2932 (86) | 2932 (93) | νCH_2 |
| 2532 | | | | 2559 (00) | 2557 (00) | 2561 (10422) | 2571 (868) | $\nu \text{N}_{15}\text{H}$ |
| 2101 | | | | | | | | $\gamma \text{NH}_3^+ + \tau \text{NH}_3^+$ |
| 1657 | 1655 (700) | 1620 (700) | $\nu_{\text{as}} \text{-COO}^-$ | 1656 (00) | 1643 (00) | 1650 (11) | 1669 (620) | $\gamma_{\text{as}} \text{NH}_3^+$ |
| 1579 | 1598 (119) | 1566 (119) | γNH_3^+ | 1562 (00) | 1565 (00) | 1562 (1443) | 1582 (3655) | $\gamma_{\text{as}} \text{NH}_3^+, \nu \text{-COO}^-$ |
| 1553 | 1581 (17) | 1550 (17) | $\beta \text{CH (ring)}, \nu \text{C}_2\text{C}_3$ | 1557 (00) | 1560 (00) | 1558 (36) | 1559 (2) | $\beta \text{CH (ring)}$ |
| 1487 | 1520 (6) | 1492 (6) | $\beta \text{CH (ring)}$ | 1479 (00) | 1471 (00) | 1474 (10) | 1475 (12) | $\beta \text{CH (ring)}, \nu \text{-COO}^-$ |
| 1455 | 1481 (46) | 1455 (46) | $\beta \text{CH (ring)}, \beta \text{H}_{21}\text{C}_{10}\text{H}_{22}$ | 1444 (00) | 1442 (00) | 1443 (170) | 1445 (179) | $\beta \text{CH (ring)}, \beta \text{CH}_2$ |
| 1409 | 1439 (357) | 1415 (357) | γNH_3^+ | 1418 (00) | 1424 (00) | 1423 (2) | 1419 (98) | $\beta \text{NH (ring)}, \beta \text{CH (ring)}$ |
| 1356 | 1399 (165) | 1377 (165) | $\gamma \text{CH}_2, \gamma \text{NH}$ | 1359 (00) | 1355 (00) | 1355 (101) | 1358 (465) | $\beta \text{NH (ring)}, \beta \text{CH (ring)}$ |
| 1339 | 1365 (185) | 1344 (185) | $\delta \text{ ring}, \nu \text{CC}$ | 1330 (00) | 1339 (00) | 1330 (214) | 1337 (10) | γCH_2 |
| 1314 | 1346 (40) | 1326 (40) | γCH_2 | 1312 (00) | 1311 (00) | 1311 (202) | 1311 (108) | γCH_2 |

| | | | | | | | | |
|------|--------------|--------------|--|--------------|--------------|---------------|---------------|--|
| 1249 | 1276 (8) | 1259 (8) | β CH (ring) | 1245 (00) | 1243 (00) | 1245 (18) | 1244 (60) | β CH (ring) |
| 1229 | 1250 (38) | 1234 (38) | β CH (ring), β NH (ring) | 1220 (00) | 1222 (00) | 1220 (138) | 1220 (6) | γ NH ₃ |
| 1209 | 1209 (10) | 1194 (10) | γ (CH ₂ , CH, NH) | 1199 (00) | 1199 (00) | 1200 (10) | 1201 (20) | γ CH ₂ , γ CH, γ NH ₃ |
| 1164 | 1174 (3) | 1160 (3) | β CH (ring) | 1155 (00) | 1164 (00) | 1151 (124) | 1163 (70) | γ NH ₃ |
| 1133 | 1130 (25) | 1118 (25) | β CH (ring), γ CH, γ NH | 1124 (00) | 1132 (00) | 1122 (6) | 1132 (11) | β CH (ring), γ CH ₂ |
| 1122 | 1107 (67) | 1095 (67) | β CH (ring), β NH (ring) | 1112 (00) | 1109 (00) | 1109 (10) | 1112 (33) | β CH (ring) |
| 1095 | 1095 (92) | 1084 (92) | γ NH ₃ , β CH(ring), β NH (ring) | 1085 (00) | 1088 (00) | 1087 (95) | 1083 (168) | β CH (ring), β NH (ring) |
| 1055 | 1069 (6) | 1059 (6) | γ (NH ₃) | 1052 (00) | 1056 (00) | 1056 (33) | 1054 (42) | γ NH ₃ , β CH (ring), β NH (ring) |
| 1006 | 1029 (22) | 1020 (22) | β CH (ring) | 1004 (00) | 1004 (00) | 1004 (10) | 1002 (114) | β CH (ring) |
| 988 | 1016 (10) | 1007 (10) | ν C ₁₁ N ₁₅ | 970 (00) | 972 (00) | 979 (13) | 976 (47) | γ NH ₃ |
| 959 | 946 (10) | 938 (10) | γ CH (ring) | 949 (00) | 953 (00) | 954 (1) | 954 (11) | β CH ₂ |
| 923 | 935 (24) | 927 (24) | β (CH ₂), γ (NH ₃) | 906 (00) | 914 (00) | 912 (11) | 919 (12) | γ CH (ring) |
| 875 | 890 (2) | 883 (2) | δ ring | 865 (00) | 866 (00) | 864 (7) | 867 (18) | δ ring |
| 863 | 865 (1) | 859 (1) | γ CH (ring) | 850 (00) | 852 (00) | 851 (144) | 850 (9) | γ CH (ring) |
| 845 | 847 (69) | 841 (69) | γ CH(ring) | 839 (00) | 839 (00) | 841 (3) | 834 (52) | γ CH (ring) |
| 835 | 829 (25) | 823 (25) | γ CH (ring), β OCO ⁻ | 821 (00) | 820 (00) | 819 (6) | 820 (31) | γ CH (ring) |
| 792 | 786 (5) | 781 (5) | Δ ring | 765 (00) | 762 (00) | 764 (1) | 762 (7) | Δ ring |
| 778 | 777 (8) | 772 (8) | δ ring, γ OCO ⁻ | 757 (00) | 758 (00) | 760 (5) | 755 (25) | Δ ring |
| 766 | 767 (17) | 762 (17) | γ OCO ⁻ | 747 (00) | 747 (00) | 749 (3) | 746 (52) | Δ ring |
| 737 | 753 (77) | 748 (77) | γ CH (ring) | 735 (00) | 746 (00) | 734 (16) | 737 (191) | γ CH (ring) |
| 715 | 749 (70) | 744 (70) | γ CH (ring) | 732 (00) | 736 (00) | 729 (321) | 729 (165) | γ CH (ring) |
| 658 | 663 (4) | 658 (4) | Δ ring, β (NH ₃) | 650 (00) | 649 (00) | 652 (5) | 649 (31) | Δ ring, γ NH ₃ |
| 625 | 636 (9) | 631 (9) | Δ ring, γ CH (ring), γ NH (ring) | 629 (00) | 623 (00) | 622 (33) | 627 (112) | γ NH |

(a) The fundamental bands observed in the FT-IR spectrum of DL-tryptophan.

(b) The calculation model used for the aqueous solution phase of DL-tryptophan ; ω (calc.): the harmonic wavenumbers calculated at B3LYP/def2-TZVP level of theory ; ω (scal.): the scaled harmonic wavenumbers derived from the calculated harmonic wavenumbers by using the correlation relation (see section 2.2.).

The symbols « ν », « β », « Δ », « δ », « τ » and « γ » represent bond stretching, angle bending, out-of-plane deformation, in-plane deformation, torsion and out-of plane angle bending vibrations, respectively.

(c) The calculation model used for DL-tryptophan in solid phase ; the calculated wavenumbers obtained for the normal modes of “a_g”, “b_g”, “a_u” and “b_u” symmetry classes are listed in separate columns.

The calculated IR intensities (in parentheses) are expressed in Km/mole.

3.2.1. The 2845-3400 cm⁻¹ spectral region

The strong band observed experimentally at 3400 cm⁻¹ (Figure 2) is assigned to the N1H26 stretching vibration in both the solid state and aqueous-solution models, it is predicted at ~3397 cm⁻¹ (solid-state model) and at 3375 cm⁻¹ (aqueous-solution model). Clearly, the solid state calculations show the best agreement. The assignment of this band agrees with a previous assignment proposed by X. Coa et al. [10].

Similarly, the band observed at 3041 cm⁻¹ is assigned to N15H stretching vibration in the solid and solution models, this band was previously observed at 3079 cm⁻¹ for L-tryptophan [10], which makes it the most sensitive to the hydrogen bond network. The shift of 38 cm⁻¹ reflects the difference in energy between the N-H...O H-bonds in L- and DL-tryptophan crystals [15]. However, the weak band observed experimentally at 2961 cm⁻¹ is assigned to the indole C-H stretching vibrations (aqueous solution model), and to the N₁₅H symmetric stretching vibration (solid state model). Finally, in both models the observed band at 2845 cm⁻¹ is assigned to the CH₂ symmetric stretching vibration. The corresponding calculated bands by using the solution phase and solid phase models are predicted at 2839 and 2930 cm⁻¹, respectively. It is remarkable that the computed wavenumber by using the solid phase model is higher by about 91 cm⁻¹ than the one calculated by using the solution phase model. This large deviation may be due to the neglect of the anharmonic effects in the calculations performed in the solid phase because all the computed wavenumbers were not scaled [36], it is important to remember here that in our study only the calculated spectrum by using the solution phase model was scaled and corrected for anharmonicity.

Previous studies performed on amino acids showed that the band observed at 2840 cm⁻¹ in the IR spectrum of L-tryptophan, which was recorded by using the dissolution-spray-deposition (DSD) technique, was assigned to the CH₂ stretching mode [10]. Also, the band observed at 2829 cm⁻¹ in the experimental IR spectrum of L-tyrosine (solid phase) was assigned to the CH₂ stretching mode [19]. These studies agree with the assignment proposed here for the band observed at 2845 cm⁻¹.

3.2.2. The 1658-2845 cm⁻¹ spectral region

As mentioned above, the region ~1700-2700 cm⁻¹ in the IR spectra of amino acids is characterized by a certain number of broad-bands. Some previous studies showed that the inclusion of solvation effects in the calculations does not lead to the prediction of absorption bands in this region [10,19-21]. Our calculated spectrum using the aqueous solution model agrees with these studies. Figure 4 shows a comparison between the calculated spectra by the solid and solution models. It is remarkable that an additional band appears at around 2561 cm⁻¹ (predicted only by the solid state model) (Figure 4b). This band is remarkably the most intense in the simulated IR spectrum ($I_{ir} = 10422$ Km/mole). This result agrees with previous studies performed on glycine and L-alanine, where the computed IR spectra in the solid phase by using the PW91 functional showed the presence of two strongest bands at around 2651 cm⁻¹ [36].

Employing the periodic DFT results, this band is assigned to the N-H stretching vibration as shown in the animation (Supplemental Material S1). The appearance of this band (Figure 4b) is not surprising, recently, A. Pawlukojc et al. [37] assigned the band observed at 2480 cm⁻¹ in the IR spectrum of L-asparagine to the N-H...O stretching vibration, F. Parker et al. assigned the observed band at ~2545 cm⁻¹ in the IR spectrum of L-cysteine to the S-H stretching vibration [38]. In addition, the computed wavenumber at 2650 cm⁻¹ (solid state calculations) for Glycine was assigned to the N-H stretching vibration [36].

On the other hand, the observed band at 2101 cm⁻¹ (Figure 2) has not correspondence in the calculated spectrum by the solid state model. In the most recent works on the amino acids, a similar band observed at 2084 cm⁻¹ in the IR spectrum of L-asparagine was assigned to a

combination mode [37], but it was not assigned in the case of L-cysteine despite its appearance in the vibrational spectra [38].

The appearance of this band may be the consequence of the extensive H-bonding networks in amino acid crystals where the NH_3^+ group forms three H-bonds, which is in general a peculiarity of amino acid crystals [11,12,17,36,39-41]. In this study, and in accordance with previous work [16,34], this band is attributed to the combination mode of out-of-plane bending and torsion of NH_3^+ group. This mode is probably due to the effect of three $\text{N-H}\cdots\text{O}$ H-bonds on the NH_3^+ group as can be seen in the animation (Supplemental Material S2).

3.2.3. The region below 1657 cm^{-1}

The strong band observed at 1657 cm^{-1} is predicted, on one side, at 1620 cm^{-1} (solution model) and assigned to the pure C=O asymmetric stretching vibration (Table III). The difference between the observed and calculated wavenumbers (37 cm^{-1}) can be explained by the effect of intra-molecular $\text{N-H}\cdots\text{O}$ H-bond as shown in (Supplemental Material S3), noting that the DL-tryptophan molecule in the solid state does not have a similar H-bond [12]. On the other side, the same observed band (1657 cm^{-1}) is predicted at $\sim 1650\text{ cm}^{-1}$ using periodic DFT and assigned to the out-of-plane bending vibration of NH_3^+ group.

The band computed by using the solution phase model and seen in Fig. 4a. at 1620 cm^{-1} has the largest intensity ($I_{\text{r}} = 700\text{ Km/mole}$), which is in general agreement with previous calculations [10]. However, a qualitative comparison of the relative intensities of the observed bands at 1657 and 1579 cm^{-1} with the calculated ones by using both the solid phase and solution phase models shows that the former model is the most suitable for interpreting these two bands.

The major contribution in the IR spectrum is due to the indole ring vibrations which are mixed with certain other vibrations as can be seen in Table III.

4. CONCLUSIONS

In this study, the structural and vibrational properties of zwitterionic DL-tryptophan were investigated using IR spectroscopy and DFT calculations performed using both the solid state and aqueous solution models.

The experimental/theoretical comparison has shown that the structural parameters and vibrational wavenumbers calculated using solid state model yield the best agreement with the experimental results. Also, the effects of the intermolecular interactions on the structural and vibrational properties of the considered compound were revealed. The results of our calculations have shown that only the calculations performed using the solid state model allowed us to attribute the broad-band observed at 2532 cm^{-1} to the N-H stretching mode.

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Appendix A. The calculated wavenumbrs in solution phase for the internal modes of DL-tryptophan.

| N° | ω_{cal} (cm ⁻¹) | |
|----|------------------------------------|--------------------|
| | Unscaled wavenumbers | Scaled wavenumbers |
| 1 | 31 | 17 |
| 2 | 49 | 36 |
| 3 | 54 | 41 |
| 4 | 82 | 70 |
| 5 | 146 | 136 |
| 6 | 209 | 200 |
| 7 | 223 | 214 |
| 8 | 236 | 228 |
| 9 | 260 | 252 |
| 10 | 301 | 294 |
| 11 | 306 | 299 |
| 12 | 382 | 376 |
| 13 | 433 | 428 |
| 14 | 444 | 439 |
| 15 | 469 | 464 |
| 16 | 492 | 487 |
| 17 | 551 | 546 |
| 18 | 593 | 588 |
| 19 | 601 | 596 |
| 20 | 636 | 631 |
| 21 | 663 | 658 |
| 22 | 749 | 744 |
| 23 | 753 | 748 |
| 24 | 767 | 762 |
| 25 | 777 | 772 |
| 26 | 786 | 781 |
| 27 | 829 | 823 |
| 28 | 847 | 841 |
| 29 | 865 | 859 |
| 30 | 890 | 883 |
| 31 | 935 | 927 |

| | | |
|----|------|-------------|
| 32 | 938 | 930 |
| 33 | 946 | 938 |
| 34 | 972 | 964 |
| 35 | 1016 | 1007 |
| 36 | 1029 | 1020 |
| 37 | 1069 | 1059 |
| 38 | 1095 | 1084 |
| 39 | 1107 | 1095 |
| 40 | 1130 | 1118 |
| 41 | 1147 | 1134 |
| 42 | 1174 | 1160 |
| 43 | 1209 | 1194 |
| 44 | 1250 | 1234 |
| 45 | 1276 | 1259 |
| 46 | 1316 | 1297 |
| 47 | 1332 | 1313 |
| 48 | 1346 | 1326 |
| 49 | 1362 | 1342 |
| 50 | 1365 | 1344 |
| 51 | 1385 | 1364 |
| 52 | 1399 | 1377 |
| 53 | 1439 | 1415 |
| 54 | 1449 | 1425 |
| 55 | 1478 | 1452 |
| 56 | 1481 | 1455 |
| 57 | 1520 | 1492 |
| 58 | 1581 | 1550 |
| 59 | 1598 | 1566 |
| 60 | 1607 | 1574 |
| 61 | 1643 | 1608 |
| 62 | 1651 | 1616 |
| 63 | 1655 | 1620 |
| 64 | 3004 | 2839 |
| 65 | 3056 | 2884 |
| 66 | 3129 | 2946 |

| | | |
|----|------|-------------|
| 76 | 3160 | 2973 |
| 86 | 3169 | 2981 |
| 69 | 3176 | 2987 |
| 70 | 3186 | 2995 |
| 71 | 3243 | 3044 |
| 72 | 3260 | 3058 |
| 73 | 3462 | 3229 |
| 74 | 3511 | 3270 |
| 75 | 3636 | 3375 |

ω_{cal} : calculated wavenumbers in (cm⁻¹).

$$\omega_{Scaled} = -0.000028 \omega_{unscaled}^2 + 1.034 \omega_{unscaled} - 14.77 \text{ (cm}^{-1}\text{)}$$

Appendix B. The calculated wavenumbers in the solid phase for the optical modes of DL-tryptophan. The 321 optical modes are classified into four subgroups according to the irreducible representation of the factor group of DL-tryptophan.

| a_g | b_g | a_u | b_u |
|----------------------|----------------------|----------------------|----------------------|
| 3389 | 3396 | 3397 | 3389 |
| 3170 | 3170 | 3170 | 3170 |
| 3118 | 3118 | 3117 | 3118 |
| 3108 | 3108 | 3108 | 3107 |
| 3098 | 3098 | 3100 | 3099 |
| 3094 | 3094 | 3095 | 3095 |
| 3045 | 3079 | 3057 | 3074 |
| 3021 | 3023 | 3023 | 3021 |
| 3006 | 3005 | 3005 | 3006 |
| 2933 | 2935 | 2965 | 2950 |
| 2930 | 2931 | 2932 | 2932 |
| 2559 | 2557 | 2561 | 2571 |
| 1656 | 1643 | 1650 | 1669 |
| 1633 | 1637 | 1619 | 1630 |
| 1603 | 1606 | 1607 | 1603 |
| 1562 | 1565 | 1562 | 1582 |
| 1557 | 1560 | 1558 | 1559 |
| 1540 | 1538 | 1533 | 1535 |
| 1532 | 1527 | 1518 | 1532 |
| 1479 | 1471 | 1474 | 1475 |
| 1444 | 1442 | 1443 | 1445 |
| 1440 | 1437 | 1438 | 1441 |
| 1418 | 1424 | 1423 | 1419 |
| 1385 | 1373 | 1382 | 1381 |
| 1359 | 1355 | 1355 | 1358 |

| | | | |
|------|------|------|------|
| 1349 | 1345 | 1347 | 1348 |
| 1330 | 1339 | 1330 | 1337 |
| 1312 | 1311 | 1311 | 1311 |
| 1294 | 1296 | 1297 | 1303 |
| 1279 | 1281 | 1280 | 1280 |
| 1245 | 1243 | 1245 | 1244 |
| 1220 | 1222 | 1220 | 1220 |
| 1199 | 1199 | 1200 | 1201 |
| 1155 | 1164 | 1151 | 1163 |
| 1137 | 1156 | 1135 | 1151 |
| 1124 | 1132 | 1122 | 1132 |
| 1112 | 1109 | 1109 | 1112 |
| 1085 | 1088 | 1087 | 1083 |
| 1052 | 1056 | 1056 | 1054 |
| 1043 | 1040 | 1040 | 1043 |
| 1004 | 1004 | 1004 | 1002 |
| 970 | 972 | 979 | 976 |
| 949 | 953 | 954 | 954 |
| 947 | 950 | 948 | 948 |
| 906 | 914 | 912 | 919 |
| 865 | 866 | 864 | 867 |
| 850 | 852 | 851 | 850 |
| 839 | 839 | 841 | 834 |
| 821 | 820 | 819 | 820 |
| 765 | 762 | 764 | 762 |
| 757 | 758 | 760 | 755 |
| 747 | 747 | 749 | 746 |
| 735 | 746 | 734 | 737 |
| 732 | 736 | 729 | 729 |
| 650 | 649 | 652 | 649 |
| 629 | 623 | 622 | 627 |
| 596 | 598 | 597 | 599 |
| 581 | 591 | 588 | 587 |
| 578 | 579 | 581 | 579 |
| 538 | 540 | 539 | 539 |
| 530 | 537 | 535 | 532 |
| 521 | 526 | 525 | 522 |
| 461 | 459 | 459 | 460 |
| 421 | 421 | 425 | 425 |
| 390 | 395 | 397 | 388 |
| 360 | 365 | 377 | 381 |
| 299 | 324 | 303 | 313 |
| 263 | 274 | 269 | 275 |
| 258 | 250 | 250 | 249 |
| 239 | 230 | 231 | 245 |

| | | | |
|-----|-----|-----|-----|
| 218 | 218 | 227 | 229 |
| 177 | 206 | 189 | 166 |
| 160 | 165 | 164 | 140 |
| 143 | 141 | 134 | 117 |
| 115 | 127 | 105 | 89 |
| 97 | 101 | 98 | 85 |
| 96 | 95 | 92 | 71 |
| 73 | 84 | 83 | 61 |
| 62 | 62 | 48 | 47 |
| 50 | 52 | 34 | |
| 38 | 29 | | |

a_u and b_u : the IR active modes.

a_g and b_g : the RAMAN active modes.

The wavenumbers are expressed in (cm^{-1}).